

# Structure of unstable nuclei in the $sd$ - $pf$ shell region by shell model with proper tensor force

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The structure of unstable nuclei around  $N = 20$  has drawn much theoretical and experimental interest associated with various exotic structures seen there such the disappearance/appearance of the magic number. Based on our extensive shell-model studies with the Monte Carlo shell model (MCSM) calculation [1], we have pointed out that the so-called shell evolution [2] ranging from stable to unstable nuclei plays an important role in the description of the exotic structures. In the case of the  $N \sim 20$  region, the shell evolution emerges as the shift of the magic number from  $N = 20$  to  $N = 16$  dominated by strongly attractive  $T = 0$  monopole interaction between  $0d_{3/2}$  and  $0d_{5/2}$  [1], which can be generalized into any strong  $T = 0$   $j_{<} - j_{>}$  monopole interactions [2].

As for the underlying microscopic origin of the shell evolution, it has been recently argued that the above mentioned property can be accounted for by the tensor force in a quantitative way with the  $\pi + \rho$  exchange potential [3]. Since its effect works also between  $j_{<}$  and  $j'_{>}$  orbits with different orbital angular momenta, it motivates us to reconstruct a new effective interaction with high predictive power for the  $sd$ - $pf$  shell, particularly for its cross-shell part. Comparing the  $T = 0$  tensor part among various interactions, we find that that of the “realistic” GXPF1 interaction [4] for the  $pf$  shell is very close to  $\pi + \rho$  while that of Millener-Kurath (MK) [5], which is often used as the cross-shell interaction, is much weaker. Together with [3], the shell evolution is expected to be quantitatively described by the shell model with an interaction with tensor force close to  $\pi + \rho$ . Indeed, we show, for instance, that the proton  $0d_{3/2}$ - $1s_{1/2}$  shell gap narrowing from  $N = 20$  to 28 is well described by such an interaction in a natural way without introducing any arbitrary shift of the monopole interaction. The tensor-induced shell evolution causes a certain narrowing of the proton  $1s_{1/2}$ - $0d_{5/2}$  shell gap at the same time, having large influence on the structure of  $^{42}\text{Si}$  which has recently attracted attention in the context of the possibility of a new magic nucleus [6]. In the present shell-model calculation, the  $^{42}\text{Si}$  nucleus is predicted to be no magic nucleus but to contain large deformed wave component driven by that effect, contrary to shell-model predictions with the MK interaction and [7].

[1] Y. Utsuno, T. Otsuka, T. Mizusaki and M. Honma, *Phys. Rev. C* 60, 054315 (1999); *ibid.* 64, 011301(R) (2001); Y. Utsuno, T. Otsuka, T. Glasmacher, T. Mizusaki and M. Honma, *ibid.* 70, 044307 (2004).

[2] T. Otsuka, R. Fujimoto, Y. Utsuno, B.A. Brown, M. Honma, and T. Mizusaki, *Phys. Rev. Lett.* 87, 082502 (2001).

[3] T. Otsuka, T. Suzuki, R. Fujimoto, H. Grawe, and Y. Akaishi, *Phys. Rev. Lett.* 95, 232502 (2005).

[4] M. Honma, T. Otsuka, B.A. Brown, and T. Mizusaki, *Phys. Rev. C* 65, 061301(R) (2002); *ibid.* 69, 034335 (2004).

[5] D.J. Millener and D. Kurath, *Nucl. Phys. A* 255, 315 (1975).

[6] J. Fridmann et al., *Nature* 435, 922 (2005).

[7] J. Retamosa, E. Caurier, F. Nowacki, and A. Poves, *Phys. Rev. C* 55, 1266 (1997); E. Caurier, F. Nowacki, and A. Poves, *Nucl. Phys. A* 742, 14 (2004).